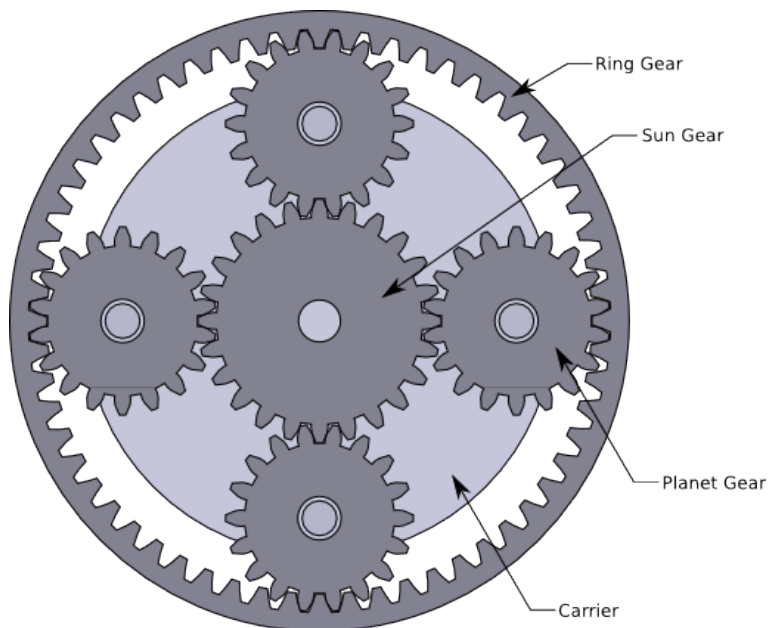


# Planetary Geartrain Analysis

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## ***Introduction***

This analysis takes an in-depth look at the governing kinematics and dynamics pertaining to a planetary geartrain. Figure 1 below shows a simplified diagram of a planetary geartrain.



*Figure 1: Simplified diagram of a planetary geartrain*

## Planetary kinematics

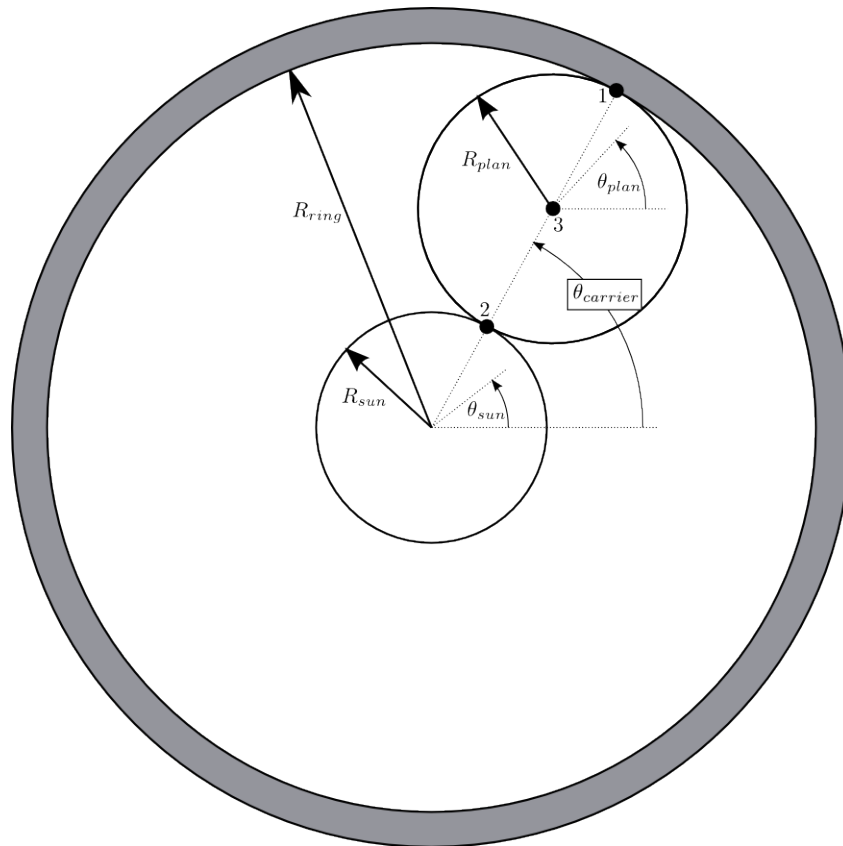


Figure 2.: Defining dimensions

- Note that point 1 on the planet gear is stationary (zero velocity) due to no-slip assumption
- Note that point 2 on the planet gear has the same velocity as point 2 on the sun gear (also due to no-slip assumption), which is given by:

$$v_{2s} = v_{2p} = R_s \dot{\theta}_s$$

- The speed of point 3 on the planet gear is  $\frac{1}{2}$  of the speed at point 2:

$$v_{3p} = \frac{1}{2} R_s \dot{\theta}_s$$

- The angular speed of the planet carrier can be found since we know the (translational) speed of the planet gear ( $V_{3p}$ ):

$$(R_s + R_p) \dot{\theta}_c = \frac{1}{2} R_s \dot{\theta}_s \Rightarrow \dot{\theta}_c = \frac{R_s}{2(R_s + R_p)} \dot{\theta}_s \quad (1)$$

- The angular speed of the planet gear can be found since we know the instantaneous speeds of two points on it (namely at points 1 and 2):

$$v_{2p} = -2R_p \dot{\theta}_p \Rightarrow \dot{\theta}_p = -\frac{R_s}{2R_p} \dot{\theta}_s \quad (2)$$

## Dynamics Analysis

### Sun Gear

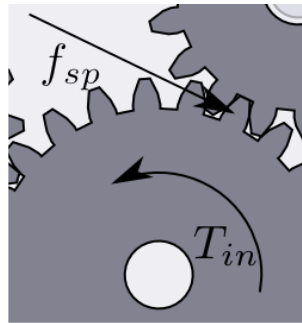


Figure 3: Sun gear free body diagram

Applying Newton's Second Law in Rotation (NSLR) to the sun gear, we have:

$$T_{in} - R_s N_p f_{sp} = J_s \ddot{\theta}_s \quad (3)$$

where  $T_{in}$  denotes the input torque (from a motor for example),  $N_p$  denotes the number of planet gears,  $f_{sp}$  denotes the tangential force at the interface between the sun gear and each planet gear (we assume here that all planet gears share the load evenly), and  $J_s$  denotes the polar moment of inertia of the sun gear.

### Planet Gear

For each planet gear, we can write NSLR as follows:

$$R_p f_{pr} - R_p f_{sp} = J_p \ddot{\theta}_p \quad (4)$$

where  $f_{pr}$  denotes the tangential force between each planet gear and the fixed ring gear and  $J_p$  denotes the polar moment of inertia of each planet gear.

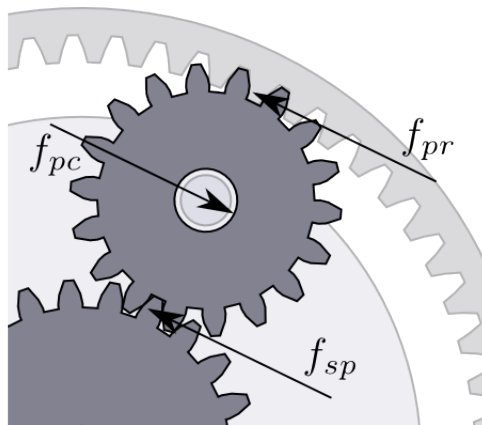


Figure 4: Planet gear free body diagram

We can also write Newton's Second Law (in translational form) in the direction of the three forces shown above (aligning the y-axis with that direction):

$$f_{sp} + f_{pr} - f_{pc} = m_p \ddot{y}_p$$

The linear acceleration,  $\ddot{y}_p$ , can be replaced by  $(R_s + R_p) \ddot{\theta}_c$  to yield:

$$f_{sp} + f_{pr} - f_{pc} = m_p (R_s + R_p) \ddot{\theta}_c \quad (5)$$

## Planet Carrier

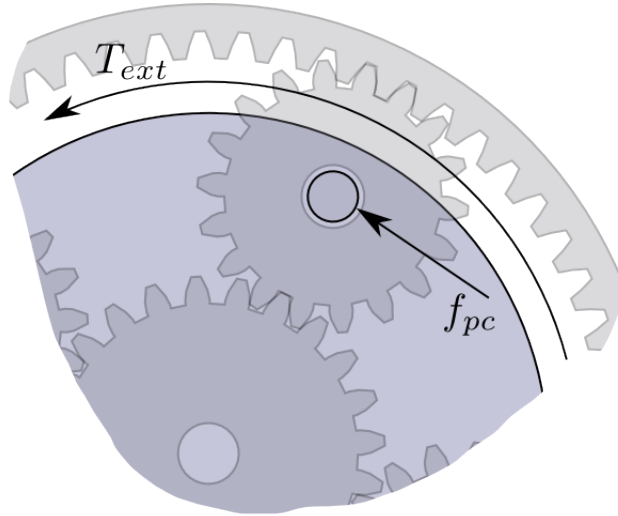


Figure 5: Planet carrier free body diagram

Applying NSLR to the planet carrier, we get:

$$T_{ext} + (R_s + R_p) N_p f_{pc} = J_c \ddot{\theta}_p \quad (6)$$

## Combining them all together

We now have 6 governing equations with 6 unknowns ( $f_{sp}$ ,  $f_{pr}$ ,  $f_{pc}$ ,  $\ddot{\theta}_p$ ,  $\ddot{\theta}_c$ , and  $\ddot{\theta}_s$ ). Through successive substitution, we can express the governing dynamical equation as follows:

$$T_{in} + \frac{R_s}{2(R_s + R_p)} T_{ext} = \left[ J_s + \left( \frac{R_s}{2R_p} \right)^2 N_p J_p + \frac{R_s^2}{4} N_p m_p + \left( \frac{R_s}{2(R_s + R_p)} \right)^2 J_c \right] \ddot{\theta}_s$$

Rearranging slightly, we can identify the contributions of the sun gear, planet gears, and planet carrier to the total effective inertia:

$$T_{in} + \frac{R_s}{2(R_s + R_p)} T_{ext} = \underbrace{\left[ \underbrace{J_s}_{\text{sun gear}} + \underbrace{N_p \frac{R_s^2}{4} \left( \frac{1}{R_p^2} J_p + m_p \right)}_{\text{planet gears}} + \underbrace{\left( \frac{R_s}{2(R_s + R_p)} \right)^2 J_c}_{\text{planet carrier}} \right]}_{\text{total effective inertia}} \ddot{\theta}_s$$